A 100-MESFET Planar Grid Oscillator

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Abstract—In this work we present a 100-MESFET oscillator which gives 21 W of CW effective radiated power (ERP) with a 16 dB directivity and a 20% dc to RF conversion efficiency at 5 GHz. The oscillator is a planar grid structure periodically loaded with transistors. The grid radiates and the devices combine quasi-optically and lock to each other. The oscillation occurs also be quasi-optically injection-locked to an external signal. The planar grid structure is very simple. All of the devices share the same bias, and they can be power and frequency tuned with a mirror behind the grid or dielectric slabs in front of it. An equivalent circuit for an infinite grid predicts the mirror frequency tuning. The planar property of the oscillator offers the possibility of a wafer-scale monolithically integrated source. Thousands of active solid-state devices can potentially be integrated in a high-power source for microwave or millimeter-wave applications.

1. INTRODUCTION

In the microwave and millimeter-wave regions solid-state oscillators have limited output power levels. However, the alternative high-power sources are tubes, which are bulky, very expensive, have a limited lifetime, and require high-voltage power supplies. Combining a large number of low-power, solid-state devices becomes attractive. Mink [1] proposed a power-combining scheme where a grid loaded with devices is placed in a quasi-optical resonator. Recently, MESFET [2], [3] and Gunn-diode [4] quasi-optical power-combining oscillators have been presented with effective radiated powers (ERPs) as high as 37 W [2]. However, these oscillators are three-dimensional structures not suitable for monolithic wafer-scale integration. In addition, the Gunn-diode array requires a separate bias supply for each diode.

In the meantime, solid-state technology has continued to improve. Recent efforts in HEMT devices [5] have resulted in a 23% power-added efficiency at 93 GHz with an output power of 18 mW. HBT's [6] have promising high-power applications with good yields and power-added efficiencies of 40% at 12 GHz with output powers of 10 W. Quantum-well devices have pushed the frequency limit to 420 GHz [7] with an output power of 200 mW. These advances in device technology make wafer-scale integrated power-combining sources very attractive. It should be possible to manufacture the grid oscillator presented here as a planar monolithic circuit which can combine a very large number of two- or three-terminal devices. This approach could give very high power solid-state oscillators in the 30–40 GHz range and reasonable power levels from low-power devices above 100 GHz. The grid oscillator is an antenna as well as the source. High directivity and a high ERP can be achieved at millimeter-wave frequencies, where a 2 in. wafer is many wavelengths across.

The planar grid oscillator is shown in Fig. 1. The transistors are Fujitsu GaAs MESFET's (FSC11LF), made for satellite receivers at 4 GHz, which typically produce 20 mW in a microstrip oscillator. The devices in adjacent rows share the gate and drain bias lines. The symmetry of the grid forces cancellation of RF currents along the horizontal bias lines, so that the radiation is primarily from the vertical drain and gate leads. The grid is made on a 0.5-mm-thick, low-dielectric-constant ($\varepsilon_r = 2.2$) material on top of a 2.5-mm-thick Duroid slab with $\varepsilon_r = 10.5$. There is a flat metal mirror behind the grid. The mirror and the grid form a Fabry–Perot cavity about a wavelength long. When the bias to the transistors is turned on, oscillations build up, and all the cavity modes compete. The mode with the lowest diffraction loss per round-trip survives, and the devices in the grid lock to this mode. This grid locks at 5 GHz. Fig. 2 shows the measured pattern with the mirror 57 mm away from the back of the substrate. The ERP is 22 W, and the directivity is 16 dB. The RF to dc efficiency is 20%, and the isotropic conversion gain [8] is 8.5 dB. The cross-polarized component is 30 dB below the copolar one at the pattern peak.

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II. EQUIVALENT CIRCUIT

In an infinite grid with all the devices locked in phase, symmetry allows us to represent the grid as an equivalent waveguide unit cell. This waveguide has magnetic walls on the sides and electric walls on the top and bottom, as shown in Fig. 3. It extends to the +z and −z directions, with the grid in the z = 0 plane. The propagating mode is TEM. The evanescent modes present reactances at the device terminals. These reactances are found by the EMF method [9]. This simple approach is justified by the fact that the period is only \( \lambda / 8 \). The reactances depend on the electrical thickness of the dielectric which supports the grid, on the geometry of the grid, and on the assumed current distribution on the transistor leads, shown in Fig. 4. The results presented here are derived from a first-order assumption for the currents: a uniform current in the gate and drain leads, and a linear current in the source lead. The latter is set to zero at the magnetic walls to satisfy the boundary conditions.

The embedding circuit for a device in the grid is shown in Fig. 5. The inductors in the gate and the drain are the lead inductances. The derivations of the expressions for the circuit elements are somewhat lengthy, but they essentially follow the procedure in Eisenhart and Kahn’s paper for a post in a waveguide [10]. The circuit elements obtained in this way are given by

\[
L = \frac{1}{j\omega} \sum_{m \text{ even}}^{\infty} \frac{\text{sinc}^2 \left( \frac{m\pi w}{2a} \right)}{Z_{\text{mol}}} Z_{\text{TE}}^{m0} \quad (1)
\]

\[
L_S = \frac{1}{j\omega} \left( \frac{2a}{a} \right)^2 \sum_{m \text{ odd}}^{n} \frac{1}{k_{mn}^2} \text{sinc}^2 \left( \frac{n\pi w}{2a} \right) \times \text{sinc}^2 \left( \frac{m\pi w}{2a} \right) \left( \frac{m}{n} + \frac{n}{m} \right) a \left( a - \frac{w}{n} \right) \left( a - \frac{w}{m} \right) Z_{\text{TE}}^{m0} \quad (2)
\]

\[
C_S = \frac{1}{j\omega} \left( \frac{2a}{a} \right)^2 \sum_{n \text{ odd}}^{n} \frac{1}{k_{mn}^2} \text{sinc}^2 \left( \frac{n\pi w}{2a} \right) \left( \frac{m}{n} \right) \left( \frac{n}{m} \right) a \left( a - \frac{w}{n} \right) \left( a - \frac{w}{m} \right) \times \text{sinc}^2 \left( \frac{m\pi w}{2a} \right) \left( 1 - \frac{a}{a - \frac{w}{n}} \right) \left( 1 - \frac{a}{a - \frac{w}{m}} \right) Z_{\text{TM}}^{m0} \quad (3)
\]

where \( k_{mn}^2 = \left(2\pi / \lambda \right)^2 - \left(\frac{m\pi}{2a}\right)^2 - \left(\frac{n\pi}{2a}\right)^2 \), and \( Z_{\text{TM}}^{m0} \)
and \( Z_{TM} \) are the impedances of the \( mn \)-th TE and TM modes [9]. Each of these is a parallel combination of the impedances in the \( +z \) and \( -z \) directions. Free space is represented as a 377 \( \Omega \) transmission line, and the mirror is represented by a shunt reactance \( jB \).

The entire grid is in this way reduced to a one-port equivalent. The measured small-signal transistor \( s \) parameters were used, and the saturation is taken into account by reducing \( S_{yy} \). The impedance looking into the oscillator from the 377 \( \Omega \) load was examined as a function of frequency. When this impedance becomes 377 \( \Omega \), an oscillation should occur.

We tested the equivalent circuit by building two grids with the same dimensions but with shorting strips instead of transistors, as shown in parts (a) and (b) of Fig. 6. The substrates are Rogers Duroid, 2.5 mm thick and 10 cm square, with a nominal dielectric constant of 10.5 and copper cladding on the back. The equivalent circuits for these two grids, shown in parts (c) and (d), use the same formulas for the inductances and capacitance as the transistor grid. The phase of the reflection coefficient of the grid was measured from 2 to 18 GHz in an anechoic chamber with a wide-band horn and an HP 8510 network analyzer. The calibration was quasi-optical. A microwave absorber at a 45° angle served as the matched load, and a copper sheet as the short. The copper sheet was translated to provide a delayed short as the third standard. The reflection coefficient phases are shown in Fig. 7. The theoretical phase calculated from the equivalent models is shown in solid line for comparison. The substrate electrical thickness which is used in the model was measured by stripping the grid pattern and repeating the reflection phase measurement. The close agreement between theory and experiment gives us confidence in the equivalent circuit model.

### III. Tuning Measurements

The position of the mirror has a strong effect on the frequency of oscillation, the grid output power, and the radiation pattern (Figs. 8 and 9). The equivalent circuit in Fig. 5 predicts frequency tuning as a function of the mirror position. In this simple transmission-line model, the oscillation frequency does not change if the mirror spacing is increased by \( \lambda/2 \), as shown in the solid lines in Fig. 8(a). The experimental results, shown with dots, agree with the theory. The mirror position affects the width of the main lobe and the size of the side lobes, as shown in Fig. 9.
theoretical pattern of the grid is calculated for a uniform array of short dipoles on an infinite dielectric, with an infinite mirror in the back [11]. The distances at which the theory predicts similar patterns are several millimeters larger than seen in the experiment.

IV. INJECTION LOCKING

In the measured results presented so far, the MESFET's were locked together with no external locking signal present. For some applications it is convenient to be able to lock the devices to an externally injected signal. For example, an electrically modulated low-power oscillator can frequency modulate the source. The injected signal can also stabilize the frequency, as shown in Fig. 10. The injection locking was performed quasi-optically, as shown in Fig. 11. The measured injection-locking bandwidth as a function of the injected power is shown in Fig. 12. It is proportional to the square root of the injected power or, equivalently, to the injected field magnitude, as expected [12].

Previous injection-locking analysis [12] has been based on the assumption that the device impedance is a function of power only, and the load impedance is a function of frequency. For the grid, the impedance is a function of both power and frequency, while the load, which is free space, is independent of power and frequency. This suggests a different approach. We use the equivalent circuit shown in Fig. 13(a). The grid is represented by a reflection coefficient \( s_1(p, f) \), which is a function of the transmitted power, \( p \), and the frequency, \( f \). Free space is represented by a load with a reflection coefficient of zero. The radiated signal is represented by a wave amplitude \( a \), where the magnitude of \( a \) is the square root of the radiated power. The injected signal is
Fig. 10. The (a) free-running and (b) injection-locked grid spectra. The center frequency is 5.2 GHz.

Fig. 11. Quasi-optical injection-locking measurement setup.

Fig. 12. The dependence of the grid locking range on the power ratio of the injected and free-running signal for a mirror distance of 55 mm. Written as \( c \), where the magnitude of \( c \) is the square root of the injected power. From the circuit model we can write the relation

\[ s_f(p, f) = a / c. \]  

Fig. 13(b) is a Rieke diagram showing \( 1/s_f \) in the complex plane. This is a convenient representation because the origin represents the free-running oscillation condition. The contours of constant power and frequency may be drawn as straight lines in a small neighborhood about the origin. As the frequency varies, the tip of the \( c/a \) arrow will trace out a semicircle, provided that the injected power remains the same and the deviation in the radiated power is small. The
frequency deviation varies symmetrically around the free-running frequency from $-\Delta f_{\text{max}}$ to $+\Delta f_{\text{max}}$, where $2\Delta f_{\text{max}}$ is the injection-locking bandwidth. The power deviation, $\Delta p$, inside the locking range is not symmetrical in general. From Fig. 13(b), the frequency and power of an injection locked-oscillator are given by

\begin{align}
\Delta f &= \Delta f_{\text{max}} \sin \phi \\
\Delta p &= \Delta p_{\text{max}} \sin (\phi + \theta) \\
\theta &= \arcsin (\Delta p_0 / \Delta p_{\text{max}})
\end{align}

where $\theta$ is the angle between the contours of constant $\Delta f$ and $\Delta p$, and $\phi$ changes from $-\pi/2$ to $+\pi+2$ inside the locking range. The quantity $\Delta p_0$ is the change in oscillator power at the free-running frequency after locking. This means that it is sufficient to measure $\Delta f_{\text{max}}$, $\Delta p_{\text{max}}$, and $\Delta p_0$ in order to predict the injection-locked oscillator behavior. Fig. 14 shows measured and calculated power variation of the injection-locked grid for two different mirror positions.

For a given frequency and injected signal level, there are actually two possible solutions (Fig. 15). However, only the higher power solution is stable, and this is what makes the $c/a$ locus in Fig. 13(b) a semicircle rather than a full circle. We can understand this by considering a perturbation $\delta a$ from the unstable solution. We will assume that $\delta a$ increases the transmitted power, so that the tip of the $c/a$ arrow moves down the $\Delta f$ contour. However, this increases the power further, so that the arrow does not stop until it reaches the stable solution.

The injection-locking bandwidth of the grid is a function of the incident angle of the injected signal. It is maximum for...
normal incidence and decreases sharply as the angle varies from the normal (Fig. 16). The pattern is very similar to the transmitting field pattern, which is indicated on the same figure. It might seem surprising that the injection-locking bandwidth has the same pattern as the transmitting grid, but this is actually a consequence of reciprocity. We can understand this by considering a general equivalent circuit for an injection-locked grid, shown in Fig. 17. The grid reflection coefficient is \( s_a \), and the source of the injection signal has a reflection coefficient \( s_s \). The scattering matrix \( S(\theta, \phi) \) takes into account the grid pattern, the free-space attenuation, and the horn pattern. It is linear and reciprocal, and we assume it to be matched. The radiated wave from the grid produces a wave in the far field. A wave \( a_x \), which feeds the horn from the injection source, produces the injected wave \( b_1 \) at the grid, so we have

\[
\begin{align*}
b_1 &= s_x a_2 \\
b_2 &= s_s a_1
\end{align*}
\]

where \( s_x(\theta, \phi) \) is effectively the transmitting grid pattern and \( s_s(\theta, \phi) \) is effectively a receiving pattern. By reciprocity, these must be equal. From Fig. 12, the injection-locking bandwidth is proportional to the injected field magnitude \( |b_1| \), which is in turn proportional to \( s_x(\theta, \phi) \) and to the transmitted wave \( |b_2| \). Fig. 16 confirms this interesting result.

V. Conclusion

A 100-MESFET oscillator is demonstrated. The oscillator is planar and suitable for monolithic integration. All the devices share the same biasing and tuning circuit, which makes this structure attractive for combining a very large number of solid-state devices on a single wafer. At higher frequencies, the wafer diameter measured in wavelengths increases, so the directivity should increase. In addition, the period of the grid is a fixed fraction of a wavelength, so the number of devices increases as well. For example, at 90 GHz, with a grid period of 0.5 mm (\( \lambda / 6 \)), a 50-mm GaAs wafer is 16 wavelengths across and would have 7500 devices. The grid should have a directivity of 30 dB and an ERP of 16 kW, assuming each device can produce 2 mW. The grid oscillator has the potential to overcome the limited power of individual solid-state millimeter-wave devices.

References


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