

4.3.7. Solution to Poisson's Equation for an Abrupt p-n Junction

Applying Gauss's law one finds that the total charge in the n-type depletion region equals minus the charge in the p-type depletion region:

$$Q_n = \epsilon_s |E(x=0)| = -Q_p \quad (4.3.45)$$

Poisson's equation can be solved separately in the n and p-type region as was done in section 3.1.1 yielding an expression for $E(x=0)$ which is almost identical to equation [3.1.4]:

$$|E(x=0)| = \frac{V_t}{L_{D,n}} \sqrt{2 \left(\exp \frac{f_n}{V_t} - \frac{f_n}{V_t} - 1 \right)} = \frac{V_t}{L_{D,n}} \sqrt{2 \left(\exp \frac{f_p}{V_t} - \frac{f_p}{V_t} - 1 \right)} \quad (4.3.46)$$

where f_n and f_p are assumed negative if the semiconductor is depleted. Their relation to the applied voltage is given by:

$$f_n + f_p = V_a - f_i \quad (4.3.47)$$

Solving the transcendental equations one finds f_n and f_p as a function of the applied voltage. In the special case of a symmetric doping profile, or $N_d = N_a$ these equations can easily be solved yielding:

$$f_n = f_p = \frac{V_a - f_i}{2} \quad (4.3.48)$$

The depletion layer widths also equal each other and are given by:

$$x_n = x_p = \left| \frac{Q_n}{qN_d} \right| = \left| \frac{Q_p}{qN_a} \right| = \left| \frac{\epsilon_s E(x=0)}{qN_d} \right| \quad (4.3.49)$$

Using the above expression for the electric field at the origin we find:

$$x_n = x_p = L_D \sqrt{2 \exp \frac{V_a - f_i}{2V_t} + \frac{f_i - V_a - V_t}{V_t}} \quad (4.3.50)$$

where L_D is the extrinsic Debye length. The relative error of the depletion layer width as obtained using the full depletion approximation equals:

$$\frac{\Delta x_n}{x_n} = \frac{\Delta x_p}{x_p} = \frac{1 - \sqrt{1 - \frac{V_t}{f_i - V_a} + \frac{V_t}{2(f_i - V_a)} \exp\left(\frac{V_a - f_i}{2V_t}\right)}}{\sqrt{1 - \frac{V_t}{f_i - V_a} + \frac{V_t}{2(f_i - V_a)} \exp\left(\frac{V_a - f_i}{2V_t}\right)}} \quad (4.3.51)$$

for $\frac{f_i - V_a}{V_t} = 1, 2, 5, 10, 20$ and 40 one finds the relative error to be 45, 23, 10, 5.1, 2.5 and 1.26 %.