

Polarization of Light

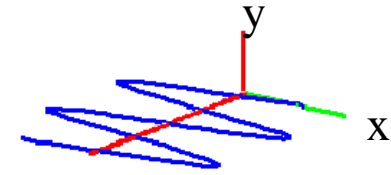
- Polarization states
- Stokes Vectors
- Mueller Matrices
- Jones Matrices
- Components
- Examples

Polarization states

Linear states of polarization

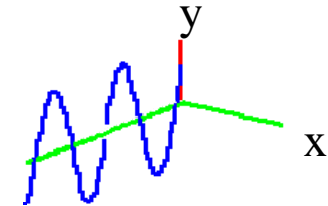
Linear x polarized

$$|E_x| \neq 0, |E_y| = 0$$



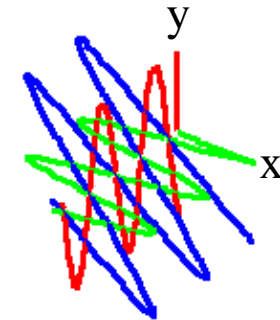
Linear y polarized

$$|E_x| = 0, |E_y| \neq 0$$



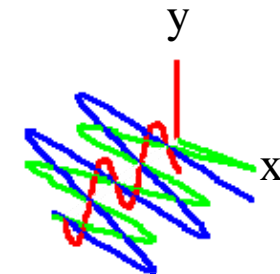
Linear -45° polarized

$$|E_x| = |E_y|, \phi_y - \phi_x = \pi$$



Linear θ polarized

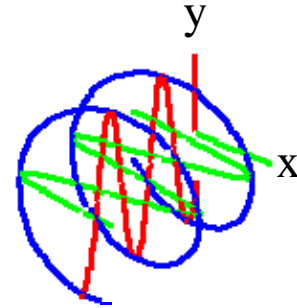
$$|E_x| \neq |E_y|, \phi_y - \phi_x = 0, \pi$$



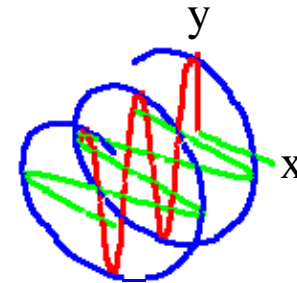
Polarization states

Elliptical states of polarization

Right-hand circular $|E_x| = |E_y|, \phi_y - \phi_x = \pi/2$



Left-hand circular $|E_x| = |E_y|, \phi_y - \phi_x = -\pi/2$

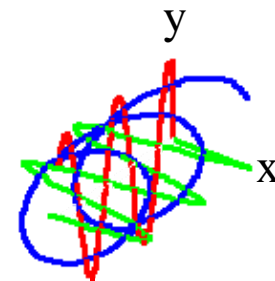
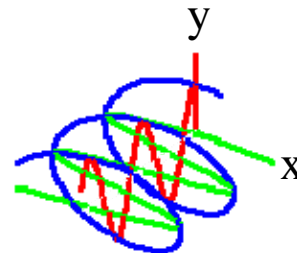


Elliptical

$$|E_x| \neq |E_y|,$$

and/or

$$\phi_y - \phi_x \neq \pm \pi/2$$



Stokes vectors

Complete description of polarization state

Perform 6 irradiance measurements with ideal polarizers:

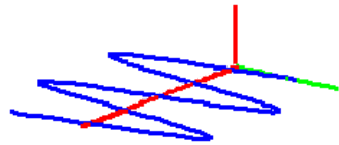
| | |
|-----------|-------------------|
| E_x | Horizontal linear |
| E_y | Vertical linear |
| E_{45} | 45° linear |
| E_{135} | 135° linear |
| E_R | Right circular |
| E_L | Left circular |

$$\vec{S} = \begin{bmatrix} S_0 \\ S_1 \\ S_2 \\ S_3 \end{bmatrix} = \begin{bmatrix} E_H + E_V \\ E_H - E_V \\ E_{45} - E_{135} \\ E_R - E_L \end{bmatrix}$$

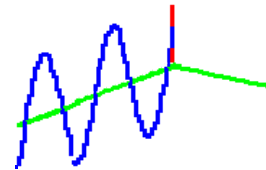
Stokes vectors

Complete description of polarization state

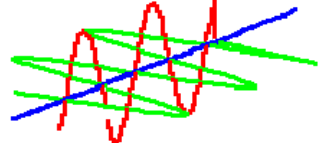
$$H = \begin{bmatrix} 1 \\ 1 \\ 0 \\ 0 \end{bmatrix}$$



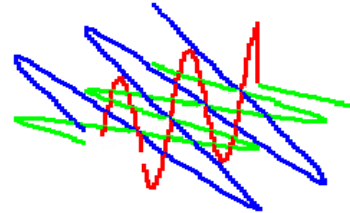
$$V = \begin{bmatrix} 1 \\ -1 \\ 0 \\ 0 \end{bmatrix}$$



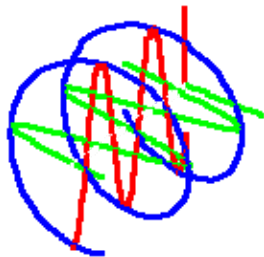
$$45 = \begin{bmatrix} 1 \\ 0 \\ 1 \\ 0 \end{bmatrix}$$



$$135 = \begin{bmatrix} 1 \\ 0 \\ -1 \\ 0 \end{bmatrix}$$



$$R = \begin{bmatrix} 1 \\ 0 \\ 0 \\ 1 \end{bmatrix}$$



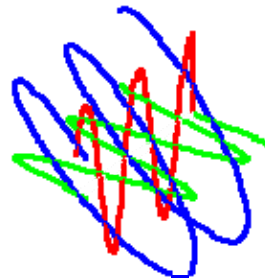
$$L = \begin{bmatrix} 1 \\ 0 \\ 0 \\ -1 \end{bmatrix}$$



$$\begin{bmatrix} 1 \\ 0 \\ .8 \\ .6 \end{bmatrix}$$

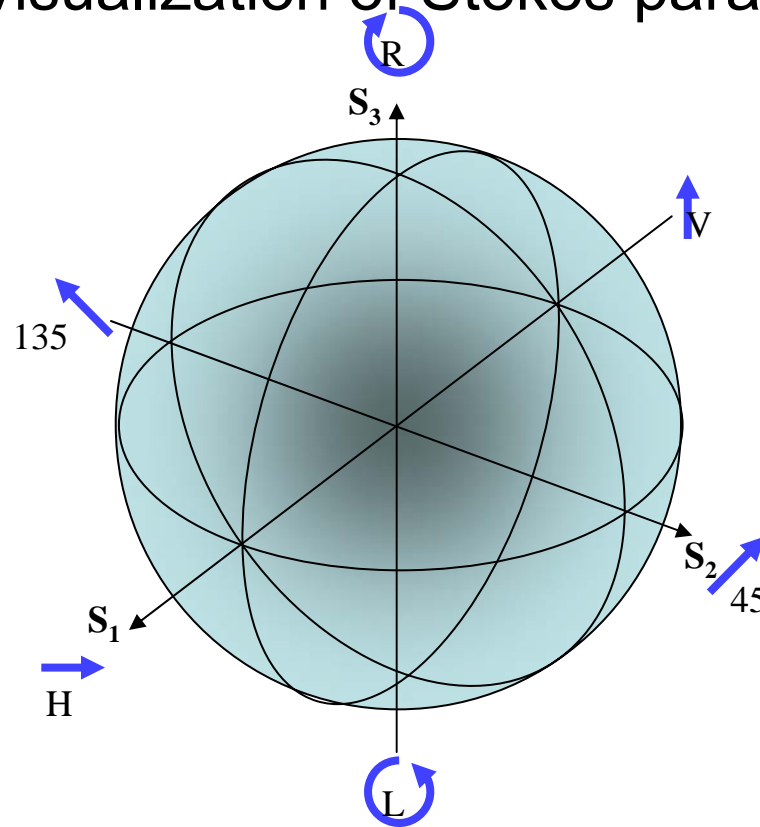


$$\begin{bmatrix} 1 \\ 0 \\ -.8 \\ -.6 \end{bmatrix}$$



Poincaré sphere

Useful visualization of Stokes parameters



$$S_i \cdot S_j = S_{i1} S_{j1} + S_{i2} S_{j2} + S_{i3} S_{j3}$$

- Surface is polarized, center is unpolarized
- Equator is linear polarized
- North hemisphere is right elliptical, south is left elliptical
- Orthogonal polarizations are on opposite points of sphere

Stokes vectors

Degree of polarization

$$DoCP = \frac{S_3}{S_0}$$

Degree of circular polarization

$$DoLP = \frac{\sqrt{S_1^2 + S_2^2}}{S_0}$$

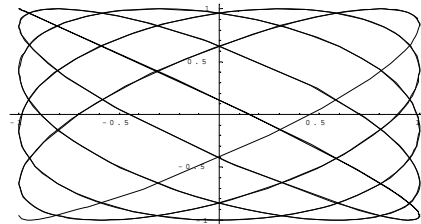
Degree of linear polarization

$$DoP = \frac{\sqrt{S_1^2 + S_2^2 + S_3^2}}{S_0}$$

Degree of polarization

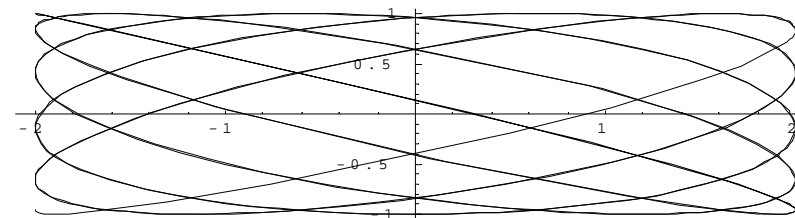
Stokes vectors are polychromatic on addition:

$$\begin{bmatrix} 1 \\ 1 \\ 0 \\ 0 \end{bmatrix} + \begin{bmatrix} 1 \\ -1 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 2 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$



H + V at different frequencies give Lissajous figure with no average polarization state

$$\begin{bmatrix} 2 \\ 2 \\ 0 \\ 0 \end{bmatrix} + \begin{bmatrix} 1 \\ -1 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 3 \\ 1 \\ 0 \\ 0 \end{bmatrix}$$



2H + V = H polarized DoP = 1/3

Mueller Matrices

Complete polarization modeling

4x4 matrix of real values describing transformation of polarizations.

Can describe de-polarizing elements.

Eigenvectors are eigenpolarization of system (unchanged).

E.g. Linear polarizer

$$\begin{bmatrix} \frac{1}{2} & \frac{1}{2} \\ \frac{1}{2} & \frac{1}{2} \\ & & & \\ & & & \end{bmatrix} \begin{bmatrix} 1 \\ 1 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 1 \\ 0 \\ 0 \end{bmatrix} \quad \begin{bmatrix} \frac{1}{2} & \frac{1}{2} \\ \frac{1}{2} & \frac{1}{2} \\ & & & \\ & & & \end{bmatrix} \begin{bmatrix} 1 \\ -1 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$

Moral: We are going to use Jones vectors/matrices to do designs, but you should know their limitations. Jones vectors deal only with systems with perfect temporal and spatial coherence. **Systems with finite coherence can be partially polarized and then you must use Mueller matrices and Stokes parameters.**

Jones vectors

Simplified for fully polarized systems

$$\vec{J} \equiv \begin{bmatrix} A_x \\ A_y \end{bmatrix}$$

where A_x and A_y are the complex amplitudes of the x and y **polarized electric fields**.

$$J_1 \cdot J_2 \equiv A_{1x} A_{2x}^* + A_{1y} A_{2y}^*$$

Inner product

$$H \equiv \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$



$$V \equiv \begin{bmatrix} -1 \\ 1 \end{bmatrix}$$

$$45 \equiv \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$



$$135 \equiv \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -1 \end{bmatrix}$$

$$R \equiv \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ j \end{bmatrix}$$



$$L \equiv \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -j \end{bmatrix}$$

Table of Jones and Mueller matrices

| Linear optical element | Jones matrix | Mueller matrix |
|---|--|--|
| Horizontal linear polarizer \leftrightarrow | $\begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$ | $\frac{1}{2} \begin{bmatrix} 1 & 1 & 0 & 0 \\ 1 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}$ |
| Vertical linear polarizer \updownarrow | $\begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}$ | $\frac{1}{2} \begin{bmatrix} 1 & -1 & 0 & 0 \\ -1 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}$ |
| Linear polarizer at $+45^\circ$ \nearrow | $\frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix}$ | $\frac{1}{2} \begin{bmatrix} 1 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}$ |
| Linear polarizer at -45° \searrow | $\frac{1}{2} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix}$ | $\frac{1}{2} \begin{bmatrix} 1 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 \\ -1 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}$ |
| Quarter-wave plate, fast axis vertical $e^{i\pi/4}$ | $\begin{bmatrix} 1 & 0 \\ 0 & -i \end{bmatrix}$ | $\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$ |
| Quarter-wave plate, fast axis horizontal $e^{i\pi/4}$ | $\begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}$ | $\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & -1 & 0 \end{bmatrix}$ |
| Homogeneous circular polarizer right \odot | $\frac{1}{2} \begin{bmatrix} 1 & i \\ -i & 1 \end{bmatrix}$ | $\frac{1}{2} \begin{bmatrix} 1 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 1 \end{bmatrix}$ |
| Homogeneous circular polarizer left \ominus | $\frac{1}{2} \begin{bmatrix} 1 & -i \\ i & 1 \end{bmatrix}$ | $\frac{1}{2} \begin{bmatrix} 1 & 0 & 0 & -1 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 1 \end{bmatrix}$ |

Rotation matrices

Components not aligned with x or y

$$\vec{J}' = R(\theta) \vec{J}$$

Transform for Jones Vectors

$$R(\theta) = \begin{bmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{bmatrix}$$

Coordinate Transform Matrix

$$T' = R(\theta) T R(-\theta)$$

Transform for Jones Matrices

Three types of physics

1. Diattenuation (polarization dependent loss)
 - Transmission is polarization dependent
 - “Polarizers”
 - A.k.a.: polarization dependent loss (PDL), dichroism
2. Retardance
 - Optical path length is polarization dependent.
 - “Wave plates”, optical activity, electro-optic
 - A.k.a.: polarization mode dispersion (PMD)
 - Poincarè sphere geometry
3. Depolarization
 - The degree of polarization may decrease depending on input polarization
 - A.k.a.: polarization scrambling

Polarizers

$$P_x \equiv \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$$

Jones matrix of linear polarizer passing H

Power transmission of analyzer and arbitrary linear polarization:

$$J_{out} = P_x L(\theta) = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} \cos \theta \\ \sin \theta \end{bmatrix} = \begin{bmatrix} \cos \theta \\ 0 \end{bmatrix}$$

$$T = J_{out} \cdot J_{out} = \cos^2 \theta$$

Malus' Law

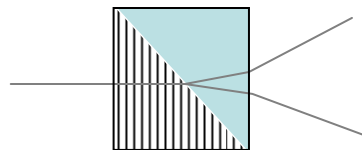
Dichroism: polarization dependent absorption

$$P_x \equiv \begin{bmatrix} e^{-\alpha z} & 0 \\ 0 & e^{-\beta z} \end{bmatrix}$$

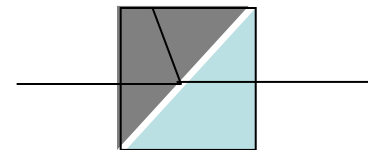
E.g.: Sheet polarizer, polarizing sunglasses, Polarcor™, wire-grid

Reflective polarizers: Brewster's angle, thin-film coatings

Crystal polarizers



E.g.: Wollaston, Rochon, Sénarmont

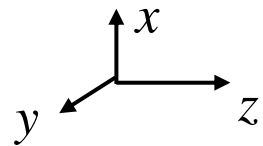


E.g.: Glan-Focault

Jones matrix example

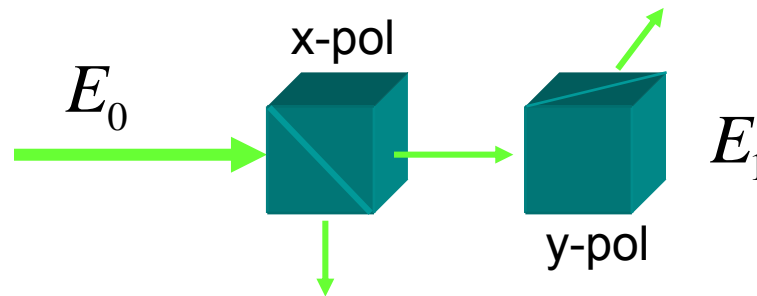
Cascaded polarizers

Crossed polarizers:



$$E_1 = \mathbf{A}_y \mathbf{A}_x E_0$$

$$\mathbf{A}_y \mathbf{A}_x = \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} = \begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix}$$

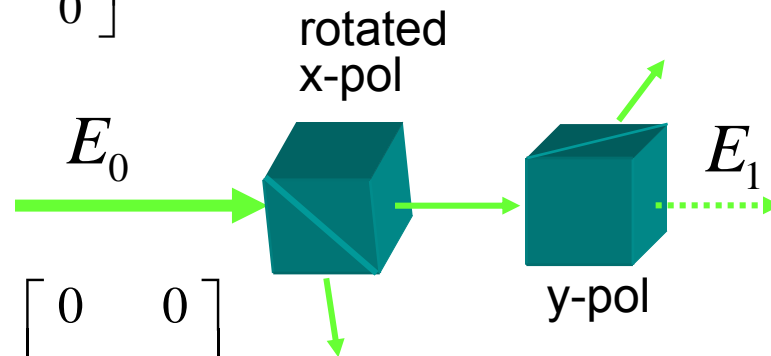


so no light leaks through.

Rotation tolerance

$$\mathbf{A}_y \mathbf{A}_x(\varepsilon) = \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & \varepsilon \\ \varepsilon & 0 \end{bmatrix} = \begin{bmatrix} 0 & 0 \\ \varepsilon & 0 \end{bmatrix}$$

$$\mathbf{A}_y \mathbf{A}_x(\varepsilon) \begin{bmatrix} E_x \\ E_y \end{bmatrix} = \begin{bmatrix} 0 & 0 \\ \varepsilon & 0 \end{bmatrix} \begin{bmatrix} E_x \\ E_y \end{bmatrix} = \begin{bmatrix} 0 \\ \varepsilon E_x \end{bmatrix} \quad \text{So } I_{out} \approx \varepsilon^2 I_{in,x}$$



Retarders

$$R \equiv \begin{bmatrix} 1 & 0 \\ 0 & e^{-j\Gamma} \end{bmatrix}$$

Jones matrix of retarder with fast axis in x

$$\Gamma = \frac{2\pi}{4} = \frac{\pi}{2}$$

$$\Gamma = \frac{2\pi}{2} = \pi$$

$$\frac{1}{n^2(\theta)} = \left\{ \frac{1}{n_o^2}, \frac{\cos^2\theta}{n_o^2} + \frac{\sin^2\theta}{n_e^2} \right\}$$

E.g. Half-wave plate.

$$\Gamma = \frac{2\pi}{\lambda_o} (n_e - n_o)L = \pi$$

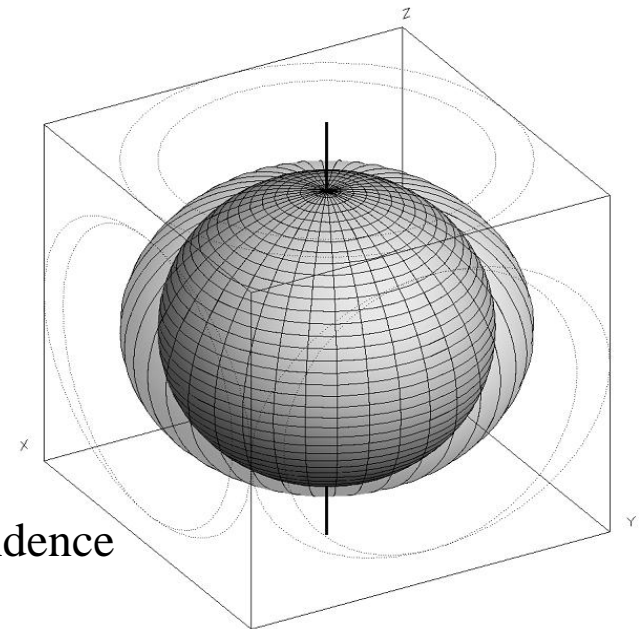
$$L = \frac{\lambda_o}{2\Delta n}$$

Quartz QWP is only 48 microns thick at 1550 nm.

Crossing plates of differing dispersion can reduce λ, θ dependence

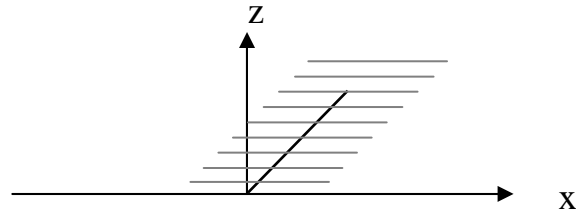
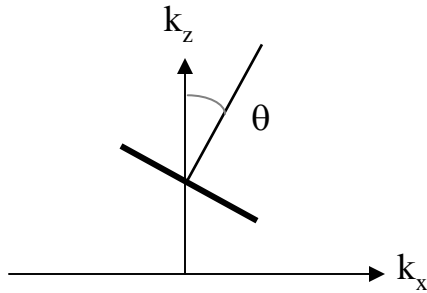
Quarter-wave plate.
Converts linear into circular.

Half-wave plate.
Converts linear into linear.



Power walk-off

Consider a Fourier propagation problem in which the k-surface is a tilted line:

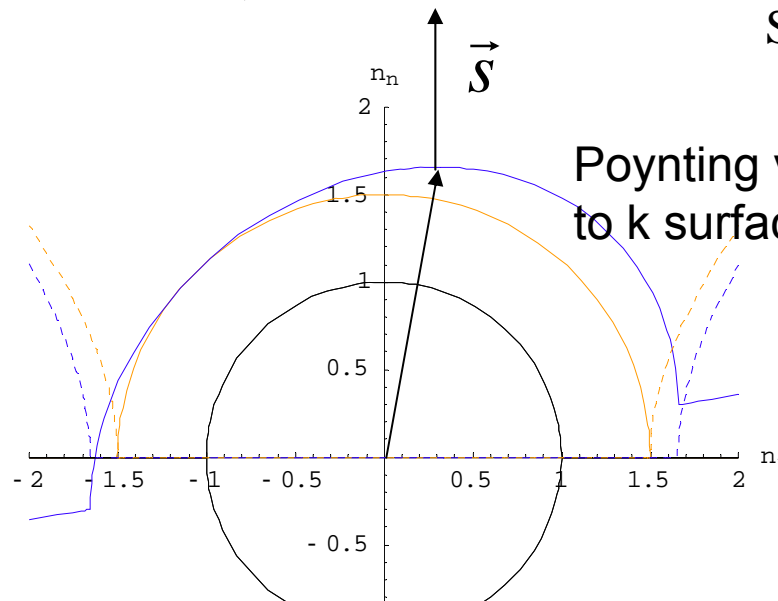


$$E(x, z) = F_x^{-1} \left\{ F_x^1 \{ E(x, 0) \} e^{-j \frac{dk_z}{dk_x} k_x z} \right\}$$

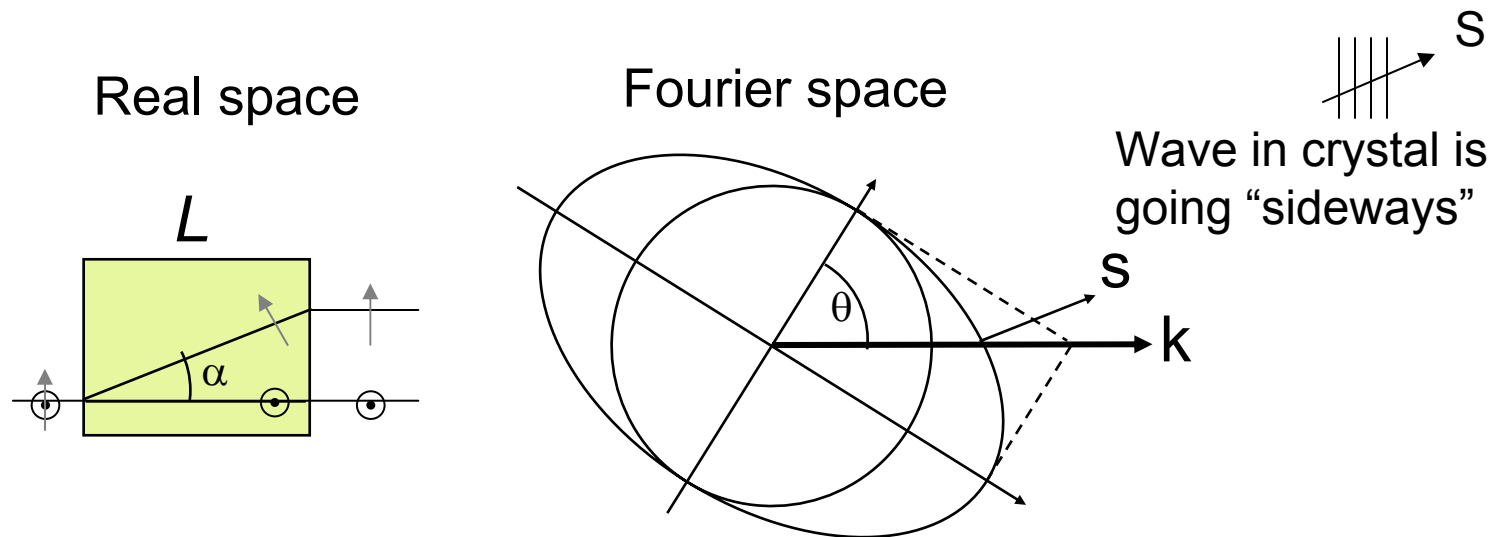
Fourier propagator

$$= E \left(x - \frac{dk_z}{dk_x} z, 0 \right)$$

Shift theorem



Beam displacing polarizer



Separates polarizations with very high isolation (limited only by crystal scatter). Emerging polarizations are parallel to a very high degree since the crystal can be flat to an arc-second.

Beam displacing polarizer

$$\alpha = \arctan \left[\frac{n_o^2}{n_e^2} \tan \theta \right] - \theta \quad \text{Walk-off angle vs. crystal cut angle}$$

$$\theta_{maxBD} = \arctan \left(\frac{n_e}{n_o} \right) \quad \text{Crystal orientation for maximum walk-off}$$

$$n_e(\theta_{maxBD}) = \left[\frac{\cos^2 \theta}{n_o^2} + \frac{\sin^2 \theta}{n_e^2} \right]^{-1/2} = \sqrt{\frac{n_o^2 + n_e^2}{2}} \quad \text{Phase velocity at this propagation angle}$$

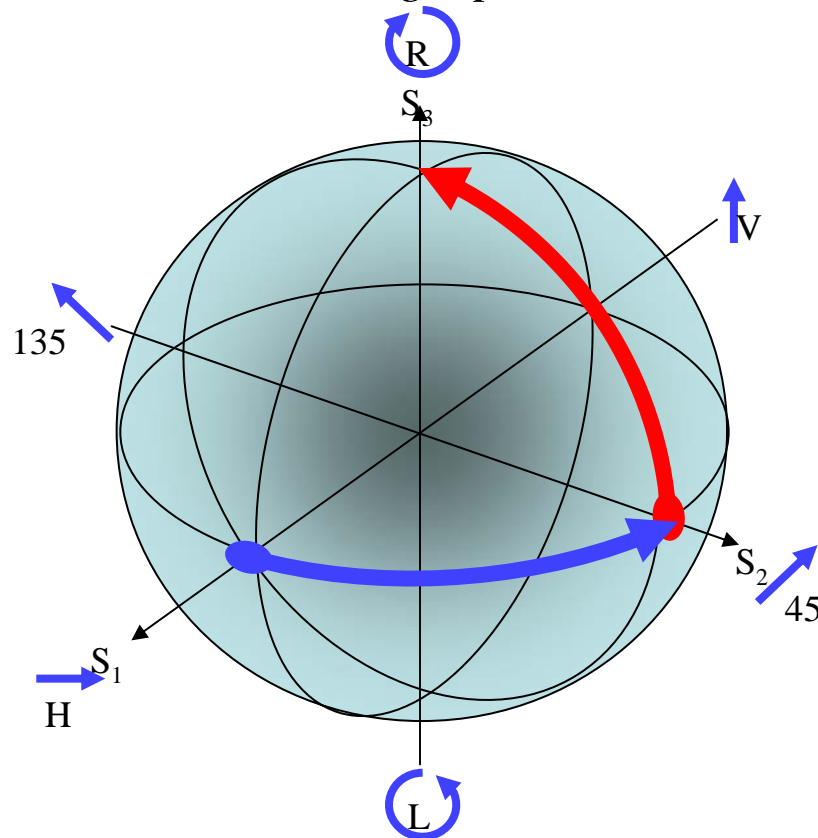
$$\Delta n = \left(\sqrt{\frac{n_o^2 + n_e^2}{2}} - n_o \right) \quad \text{Differential optical path length}$$

$$= \left(\sqrt{\frac{1.9447^2 + 2.1486^2}{2}} - 1.9447 \right) = 0.1045$$

YVO₄ at 1.55 μm

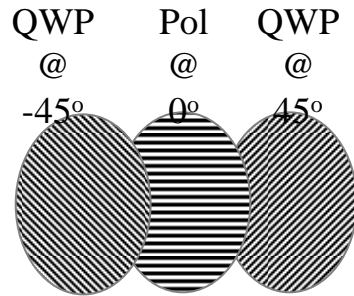
Poincarè and retarders

- Retarders rotate the polarization state on the Poincarè sphere.
- Axis of rotation connects *eigenpolarizations* of the retarder.

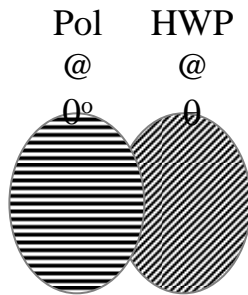


- **Example1: QWP with horizontal axis converts 45° linear into RHC**
- **Example2: Optically active (or Faraday) rotator converts H to 45°**

Useful combinations



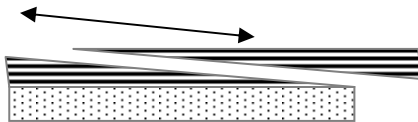
Circular polarizer



Adjustable linear polarizer



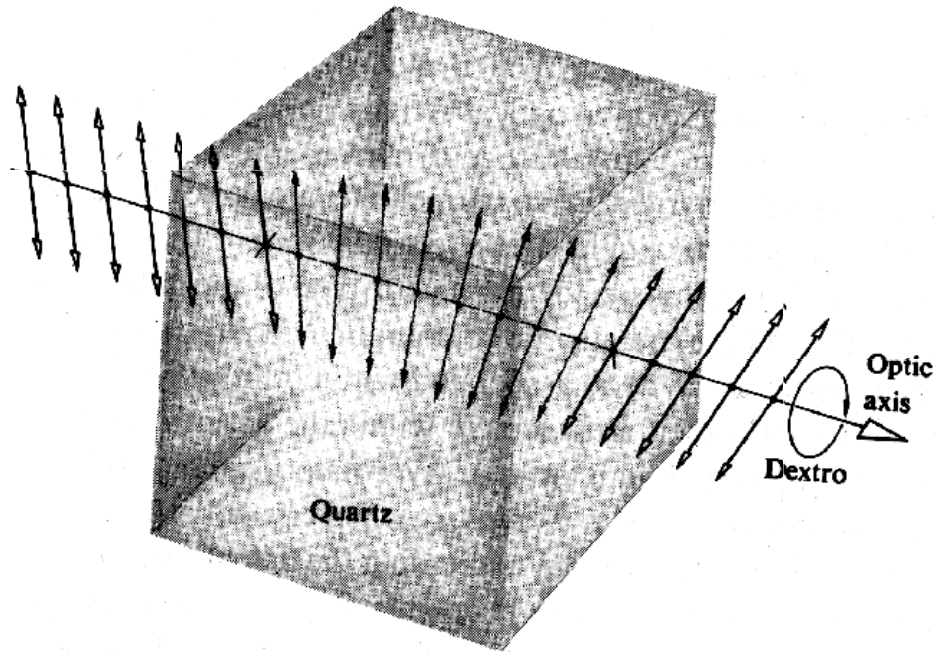
Adjustable transmission



Babinet-Soleil adjustable retarder

Optical activity

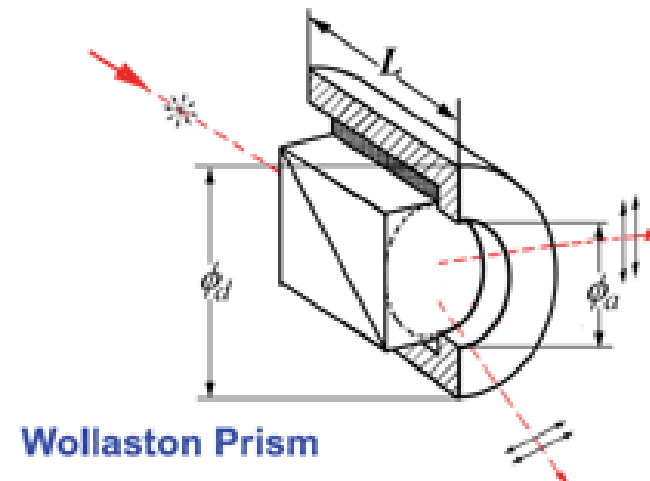
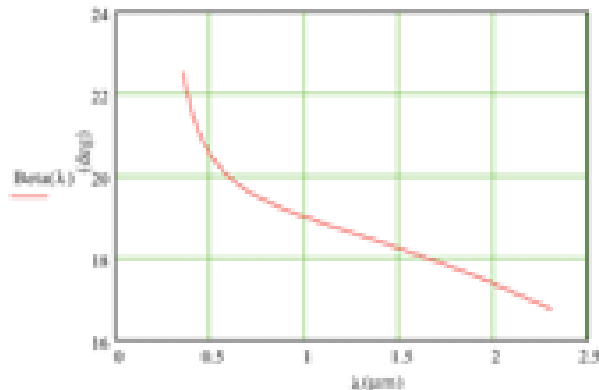
Retarders with circular eigenstates



$$OA \equiv \begin{bmatrix} \cos \delta/2 & \sin \delta/2 \\ -\sin \delta/2 & \cos \delta/2 \end{bmatrix}$$

Wollaston

- Wollaston polarizer is made of two birefringent material prisms that are cemented together. The deviations of the ordinary and extraordinary beams are nearly symmetrical about the input beam axis. The separation angle exhibits chromatic dispersion, as shown in the blow. Any separation angle can be designed upon the requirement. Typical separation angle vs wavelength is shown in the plot below for Calcite.



Extinction Ratio: 5×10^{-6}

Beam Separation Angle: 16.7-22.5 at 980nm

Faraday Rotator

- A **Faraday rotator** is an optical device that rotates the polarization of light due to the Faraday effect, which in turn is based on a magneto-optic effect.
- The Faraday rotator works because one polarization of the input light is in ferromagnetic resonance with the material which causes its phase velocity to be higher than the other.
- The plane of linearly polarized light is rotated when a magnetic field is applied parallel to the propagation direction (permanent magnets). The empirical angle of rotation is given by:

$$\beta = VBd$$

- Where β is the angle of rotation (in radians). B is the magnetic flux density in the direction of propagation (in teslas). d is the length of the path (in metres) where the light and magnetic field interact. Then V is the Verdet constant for the material. This empirical proportionality constant (in units of radians per tesla per metre, $\text{rad}/(\text{T}\cdot\text{m})$) varies with wavelength and temperature and is tabulated for various materials.
- Note that direction of rotation given by sign of B Rotation is **independent of propagation** direction (NOT like waveplate)
- Faraday rotators are used in optical isolators to prevent undesired back propagation of light from disrupting or damaging an optical system.

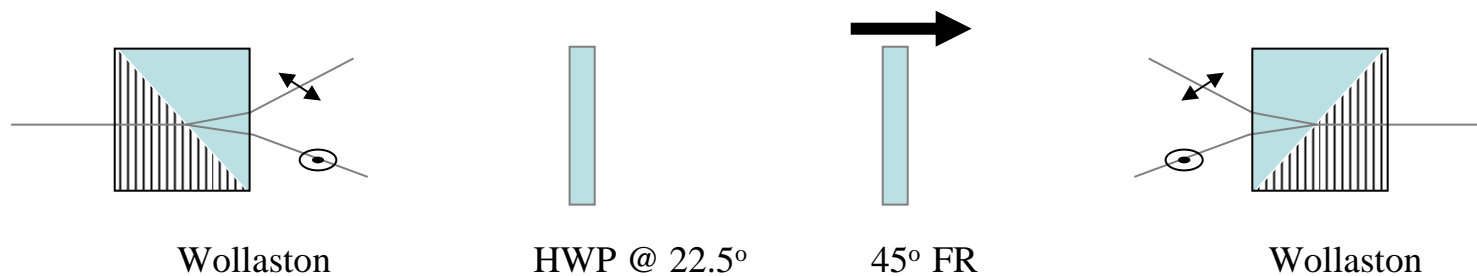
Polarization Dependent Isolator

- The polarization dependent isolator, or Faraday isolator, is made of three parts, an input polarizer (polarized vertically), a Faraday rotator, and an output polarizer, called an analyzer (polarized at 45 degrees)
- Light traveling in the forward direction becomes polarized vertically by the input polarizer. The Faraday rotator will rotate the polarization by 45 degrees. The analyser then enables the light to be transmitted through the isolator.
- Light traveling in the backward direction becomes polarized at 45 degrees by the analyzer. The Faraday rotator will again rotate the polarization by 45 degrees. This means the light is polarized horizontally (the rotation is insensitive to direction of propagation). Since the polarizer is vertically aligned, the light will be extinguished

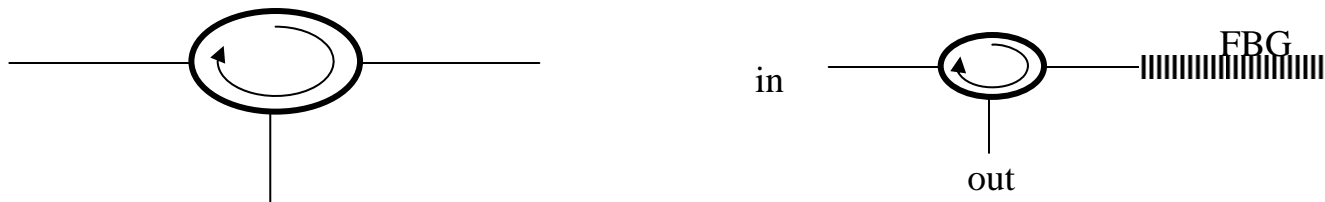


Nonreciprocal optics aka Faraday rotators

Example: Polarization **independent** isolator

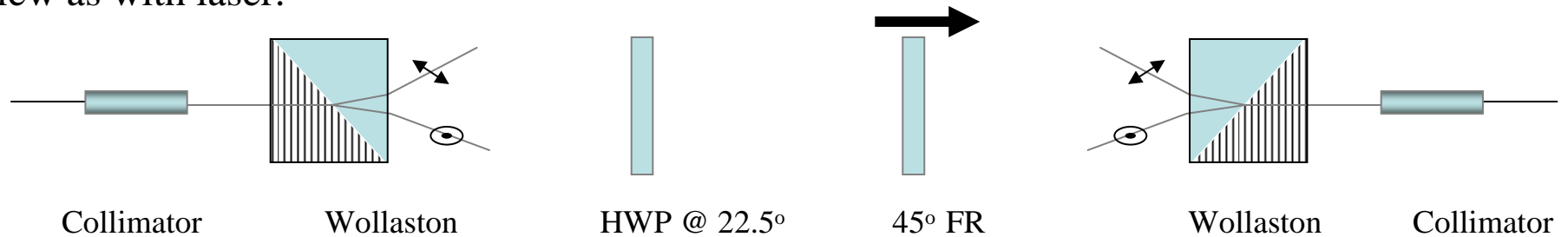


Similar optics can form a pseudo-circulator:



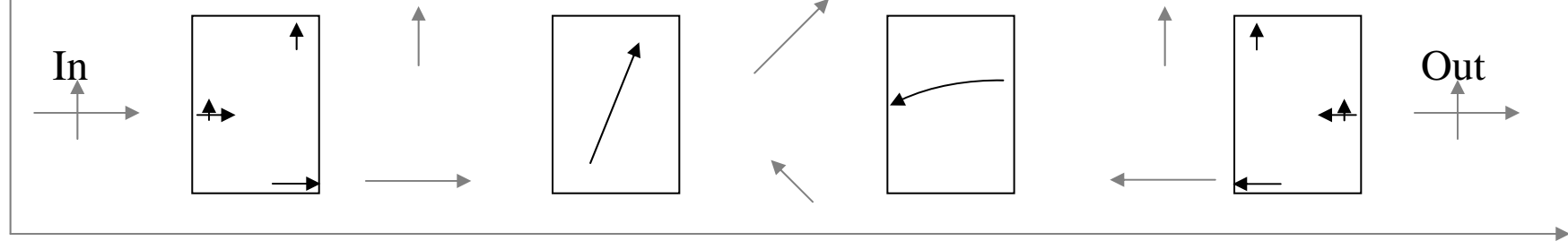
Design notation for polarization manipulation

View of polarization as if looking into laser. This is the EE convention and yields RHC with thumb pointing into viewers eyes, along with laser. To use the Physics convention, consider view as with laser.



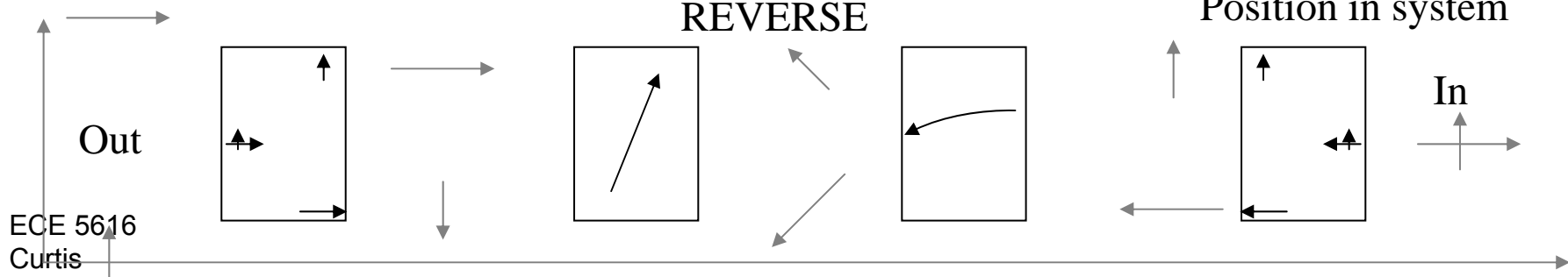
Angle of propagation

FORWARD

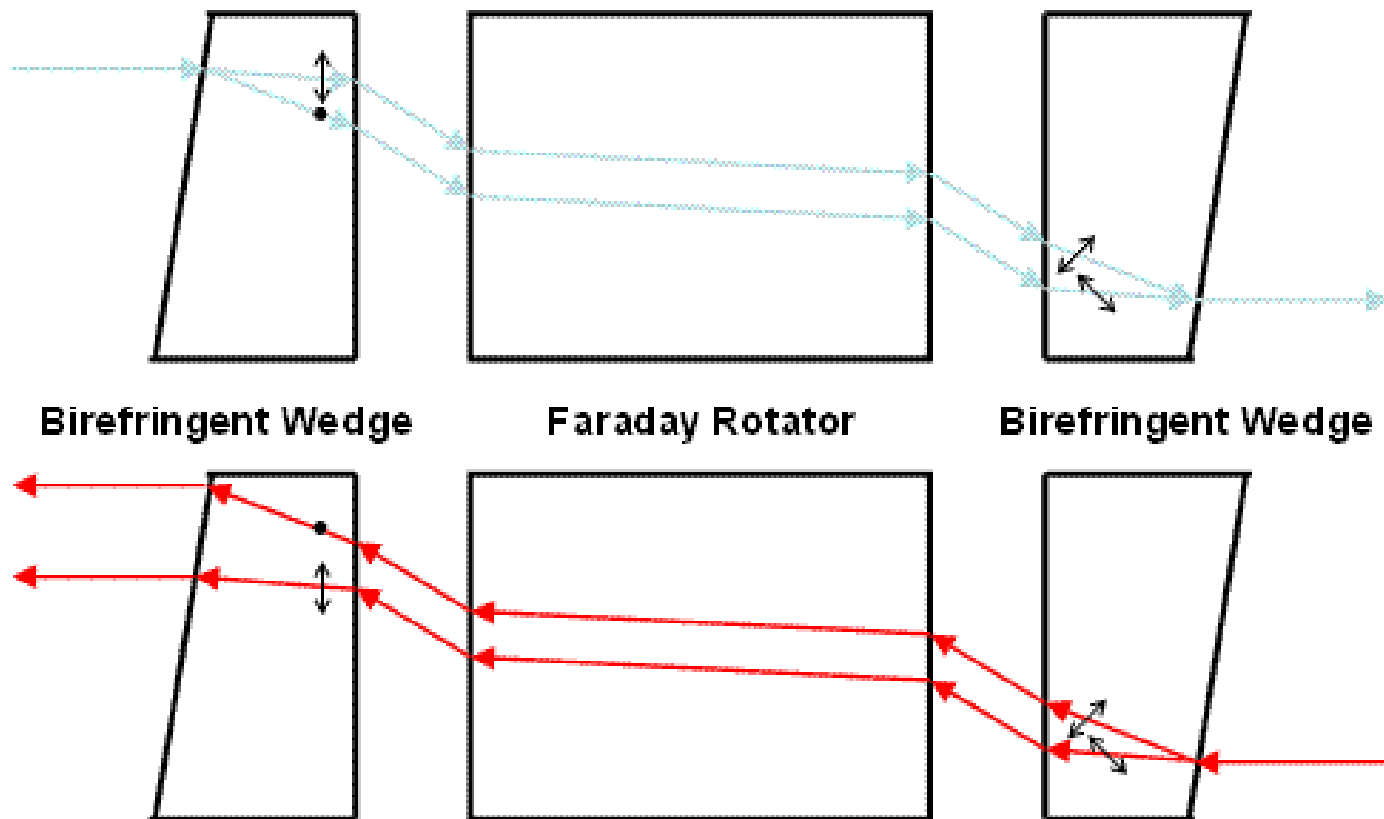


REVERSE

Position in system

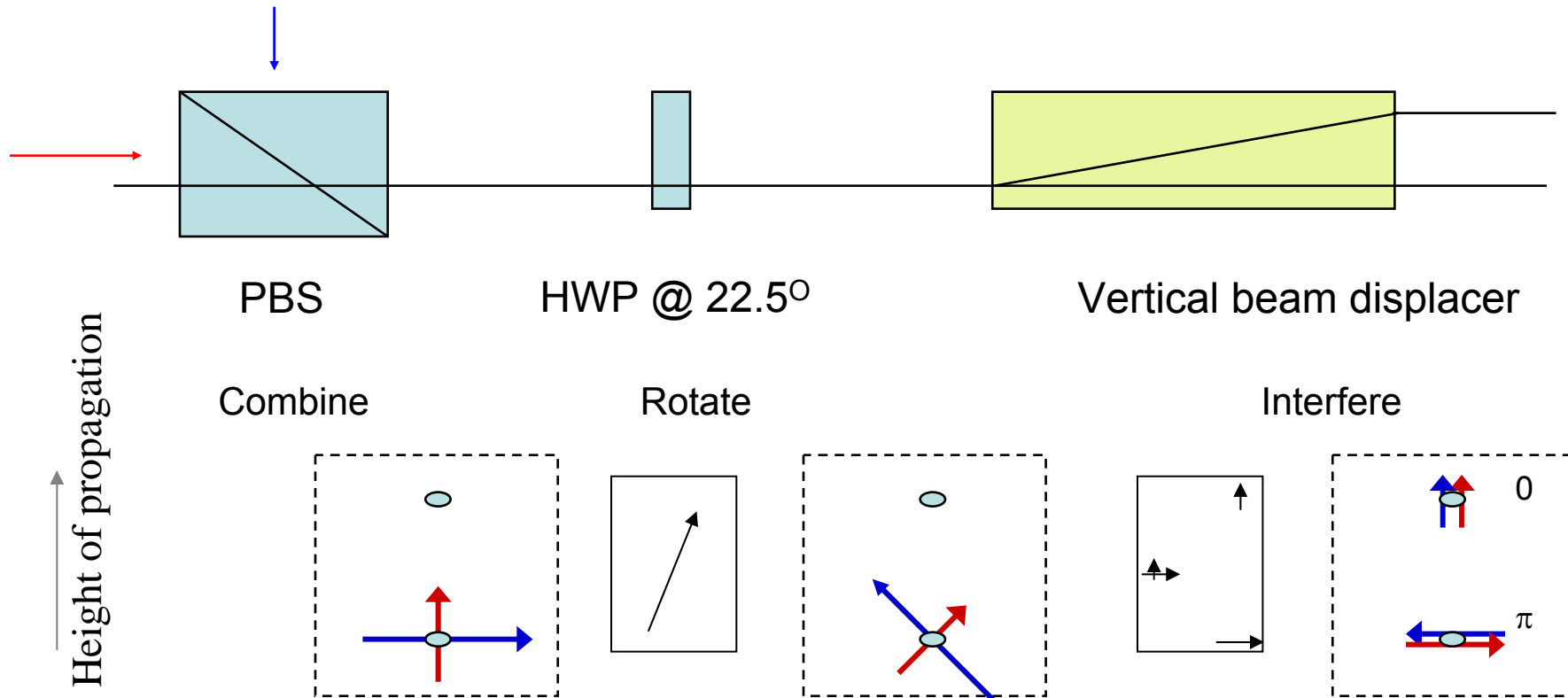


Polarization Independent Isolator

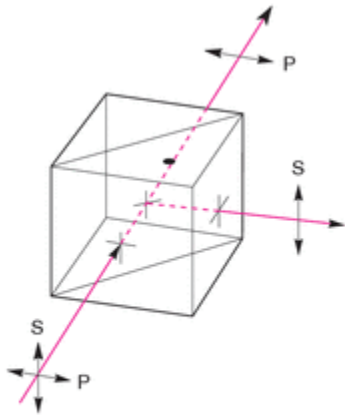


Another example

Spatial phase-shifting interferometer



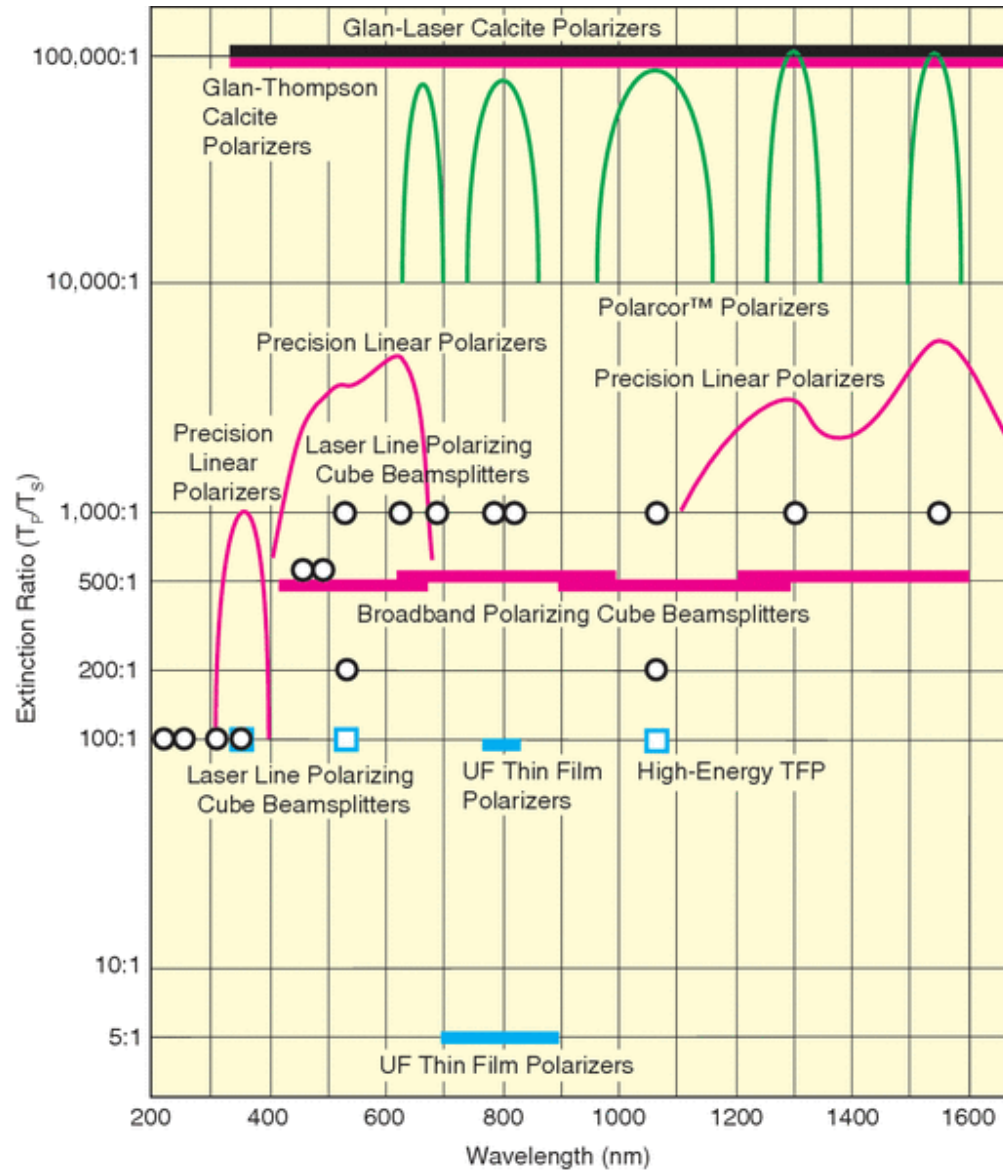
PBS Example



NOTE: To avoid damage, beam must enter prism on the side marked with a dot.

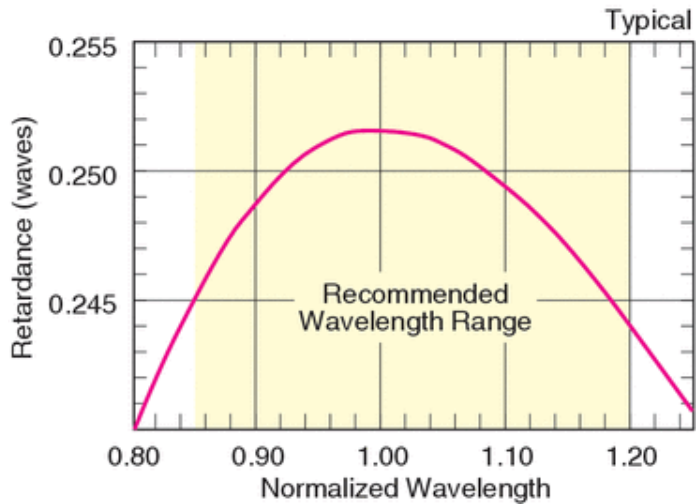
| | |
|----------------------------|--|
| Material | SF 2, NSSK grade, precision annealed optical glass |
| Wavefront Distortion | $\leq \lambda/4$ at 632.8 nm over the clear aperture |
| Clear Aperture | Central diameter, >80% of dimension |
| Surface Quality | 20-10 scratch-dig |
| Efficiency | $T_p > 80\%$, >90% average, $R_s > 99.5\%$ average |
| Extinction Ratio | $T_p/T_s > 500:1$, 1000:1 average |
| Transmitted Beam Deviation | ≤ 5 arc min |
| Reflected Beam Deviation | $90^\circ \pm 5$ arc min |
| Angle of Incidence | $0^\circ \pm 5^\circ$ |
| Dimensions Tolerance | ± 0.25 mm |
| Antireflection Coating | Broadband, multilayer coating, $R_{avg} < 1.0\%$ per surface |
| Temperature Range | $-50^\circ\text{C} - 90^\circ\text{C}$ |
| Durability | MIL-C-675C |
| Cleaning | Non-abrasive method, acetone or isopropyl alcohol on lens tissue recommended (see Care and Cleaning of Optics) Cemented optic, do not immerse in a solvent |
| Damage Threshold | 2 kW/cm ² CW, 1 J/cm ² with 10 nsec pulses, typical |

Polarizers



Newport Catalogue

Zero Order Waveplates



For $\Gamma = \lambda/2$ the plate is called a „0“ order half-wave plate. Thickness of the plate is:

$$d = \frac{\lambda}{2(n_e - n_o)}$$

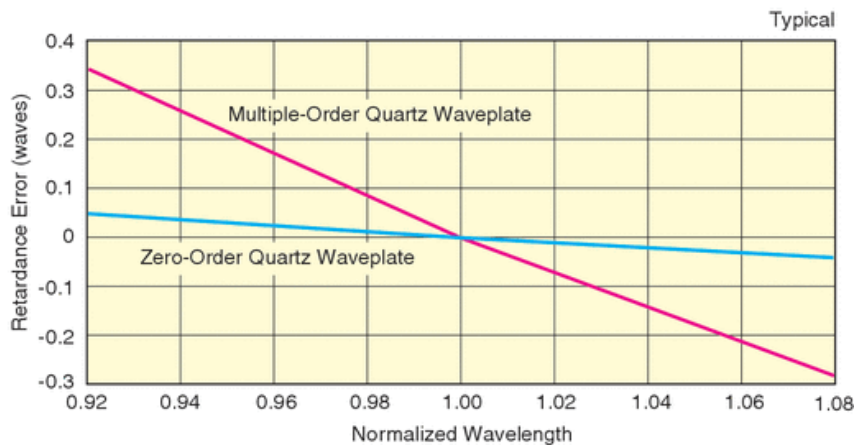


| | |
|----------------------------|--|
| Retarder Material | Birefringent polymer film stack |
| Substrate Material | BK 7, grade A, fine annealed optical glass |
| Retardation | $\lambda/4$ or $\lambda/2$ |
| Retardation Accuracy | $\pm\lambda/100$ |
| Wavefront Distortion | $\leq\lambda/4$ at 632.8 nm over the full aperture |
| Clear Aperture | 10.2 mm |
| Surface Quality | 40-20 scratch-dig |
| Transmitted Beam Deviation | ≤ 1 arc min |
| Acceptance Angle | $\pm 7^\circ$ |
| Thickness | 3.56 mm |
| Housing Diameter | 25.4 ± 0.13 mm |
| Housing Thickness | 6.2 mm |
| Temperature Range | -20 °C to 50 °C |
| Antireflection Coating | Broadband, multilayer coating, $R_{avg} < 0.5\%$ |
| Cleaning | Non-abrasive method, acetone or isopropyl alcohol |
| Damage Threshold | 500 W/cm ² CW, 0.3 J/cm ² with 10 nsec pulses, visible |
| Housing | Black anodized aluminum |

Multiple order Waveplates



Must order for laser line
 Incident angles must be close to normal
 Multiple nulls with polarizer



| | |
|--------------------------|---|
| Material | Quartz, schlieren grade |
| Construction | Single plate |
| Retardation | $\lambda/4$ or $\lambda/2$ |
| Retardation Accuracy | $\pm\lambda/300$ at $20^\circ\text{C} \pm 1^\circ\text{C}$ |
| Wavefront Distortion | $\leq\lambda/10$ at 632.8 nm over the full aperture |
| Clear Aperture | \geq central 90% of diameter |
| Surface Quality | 10-5 scratch-dig |
| Wedge | <0.5 arc sec |
| Diameter Tolerance | +0/-0.25 mm |
| Reflectivity per Surface | Single wavelength: $R < 0.5\%$ total, or 0.25% per surface Dual wavelength: $R_{\text{avg}} < 1.5\%$ |
| Thickness | Single wavelength: 1 mm, nominal Dual wavelength: 0.5–4 mm, nominal |
| Cleaning | Non-abrasive method, acetone or isopropyl alcohol on lens tissue recommended, caution: fragile, thin optic |
| Damage Threshold | 2 MW/cm ² CW, 2 J/cm ² with 10 nsec pulses |